

Indirect carrier leakage in short-wavelength InAs/AlSb quantum cascade lasers

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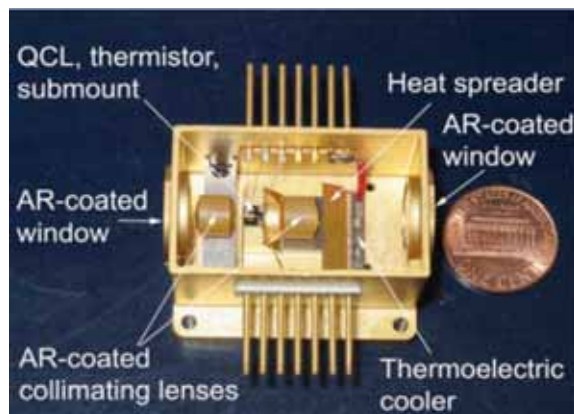
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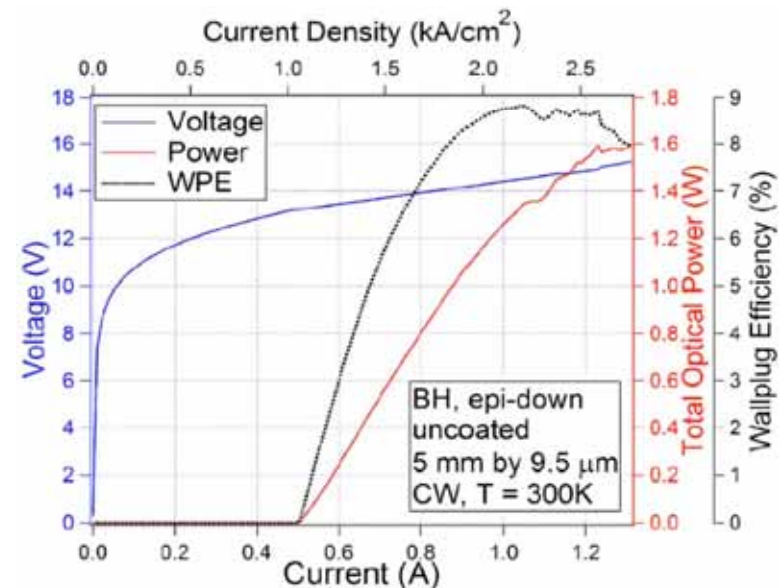
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Mid-infrared semiconductor lasers for 2-4 μm spectral region

- Applications: gas sensing, free space optical communications, etc.
- Interband QW lasers \rightarrow limitations for $\lambda > 2.5\mu\text{m}$ due to Auger recombination
- Short-wavelength QCLs (high band offsets needed $>400\text{meV}$ to provide lasing at $\lambda < 3\mu\text{m}$) – no problem of Auger recombination
- Excellent performance of QCLs at $\lambda > 4\mu\text{m}$ (InAlGaAs/InP QCL – 1.6W CW at 300K) \rightarrow efforts to reduce λ by using new material systems

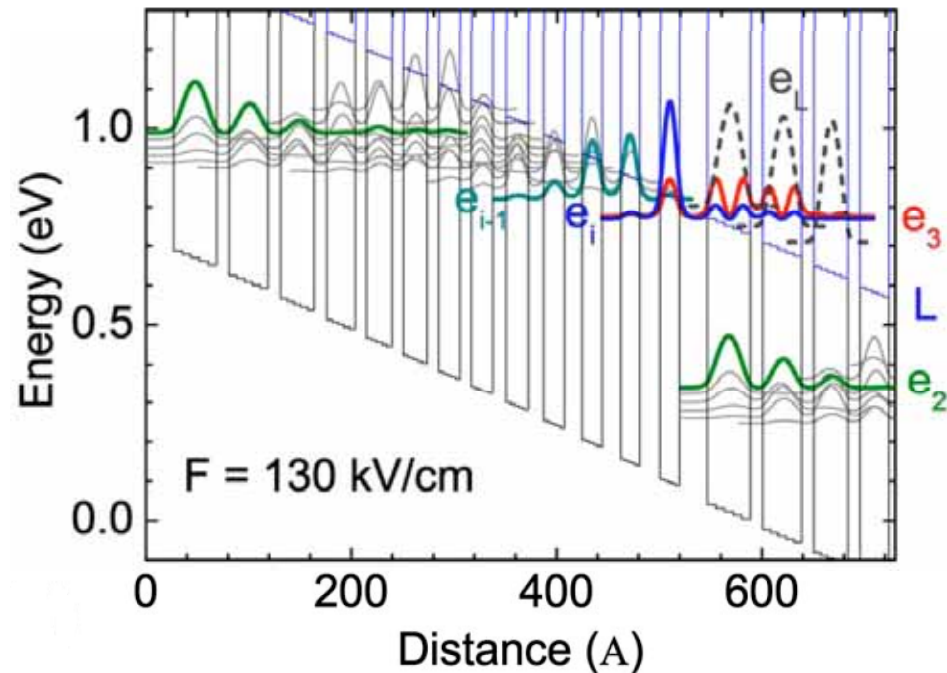


A. Lyakh et al, Appl. Phys. Lett. **92**, 111110 (2008).



Short-wavelength QCLs

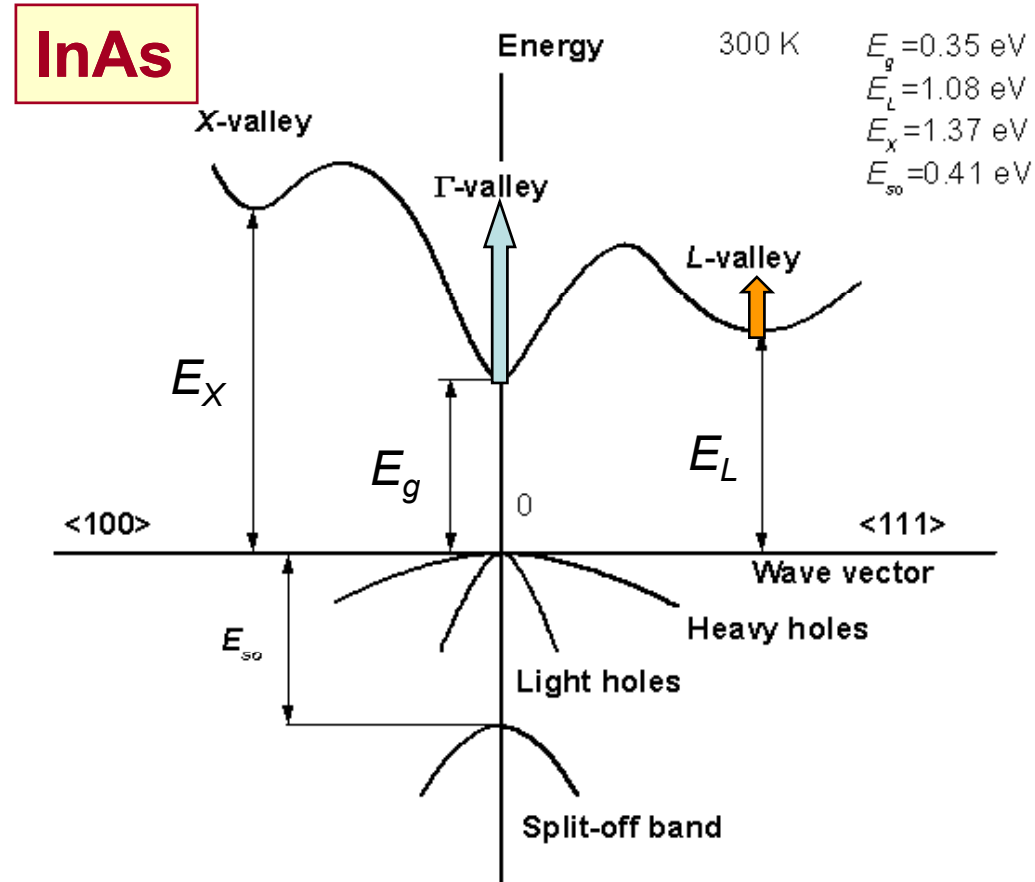
J. Devenson et al, Appl. Phys. Lett. **91**, 251102 (2007)



Conduction band diagram of InAs/AlSb 2.9 μm QCL
Solid curves - the moduli squared of the relevant electron wave functions.
Dotted lines - e_L levels associated with the L valley calculated using Γ -L separation value of 0.73 eV.

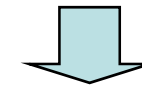
- InAs/AlSb exhibits highest conduction band Γ - Γ offset of 2.1 eV.
- Largest separation between direct and indirect minima of conduction band (Γ -L separation is 0.72-0.76 eV in InAs).
- Modelling \rightarrow Γ -L electron leakage from the upper laser levels may influence the laser performance.
- Use high hydrostatic pressure to investigate presence and importance of carrier leakage into L-valley

Effect of high-pressure on conduction band of InAs



Use hydrostatic pressure to manipulate band structure at constant T:

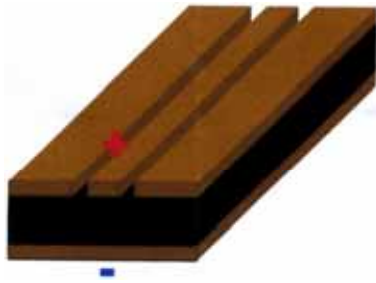
- $dE_{\Gamma}/dp = 10.2$ meV/kbar
- $dE_L/dp = 3.2$ meV/kbar
- $d(E_L - E_{\Gamma})/dp = -7$ meV/kbar



This provides a possibility to investigate the importance of carrier leakage into L-valley using high pressure

Experiment

Samples



MBE grown

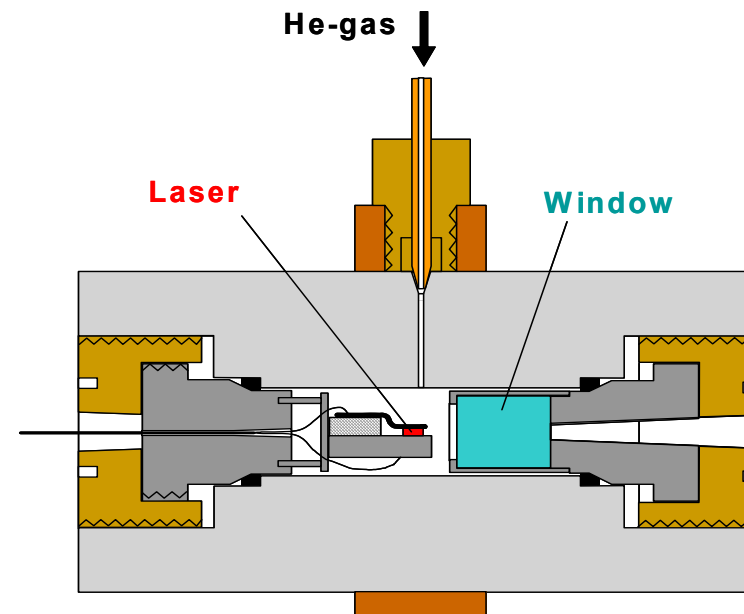
InAs/AlSb QCLs on (100) InAs substrate

$\lambda \sim 2.9\mu\text{m}$ and $3.3\mu\text{m}$ at RT

Ridge lasers (12 and 17 μm wide)

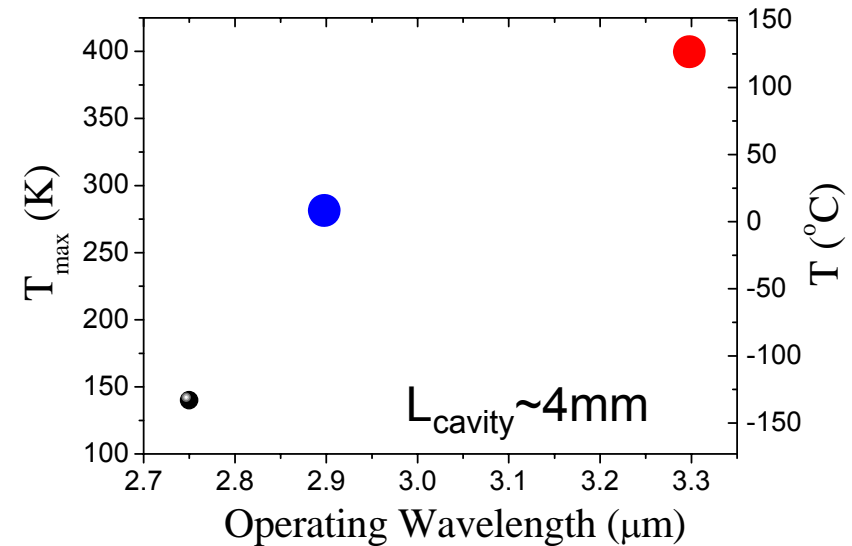
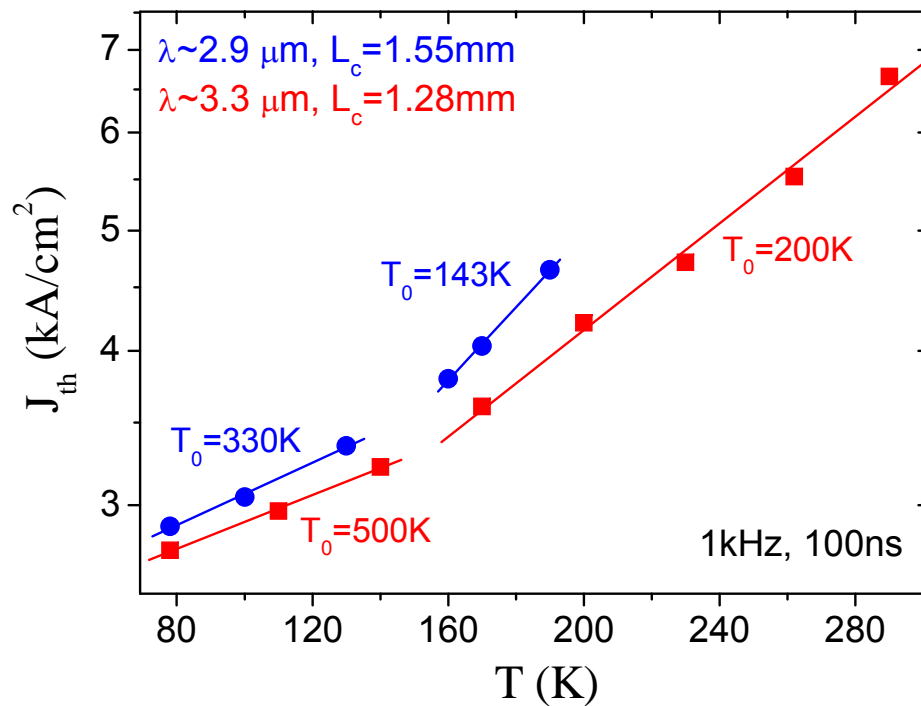
$L_{\text{cavity}} = 1.1\text{-}2.05\text{ mm}$

High-pressure cell



- Pressures up to 10kbar
- Temperatures down to 80K

Temperature dependence of I_{th}



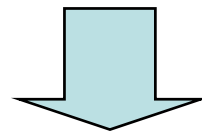
J. Devenson et al, APL **91**, 251102 (2007)

- Lower J_{th} , higher T_0 and maximum operating temperature, T_{max} , for longer wavelength devices
- Temperature sensitive loss process is more important at shorter wavelengths
- This is consistent with indirect Γ - L carrier leakage since the shorter wavelength devices have a smaller Γ - L separation

Carrier leakage as function of T and p

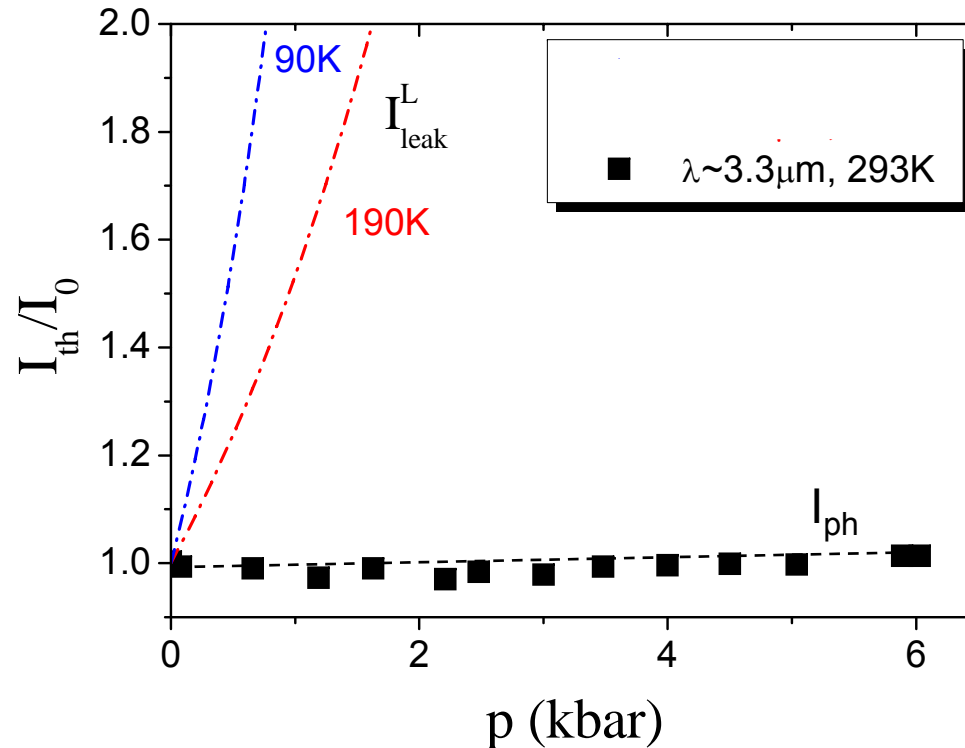
$$I_{leak}^L(T) = I_{0leak} \exp\left(-\frac{E_a}{kT}\right)$$

$$I_{leak}^L(p) \propto \exp\left(-\frac{dE_a}{dp} \frac{p}{kT}\right) \propto \exp\left(-\frac{dE_L - dE_\Gamma}{dp} \frac{p}{kT}\right)$$



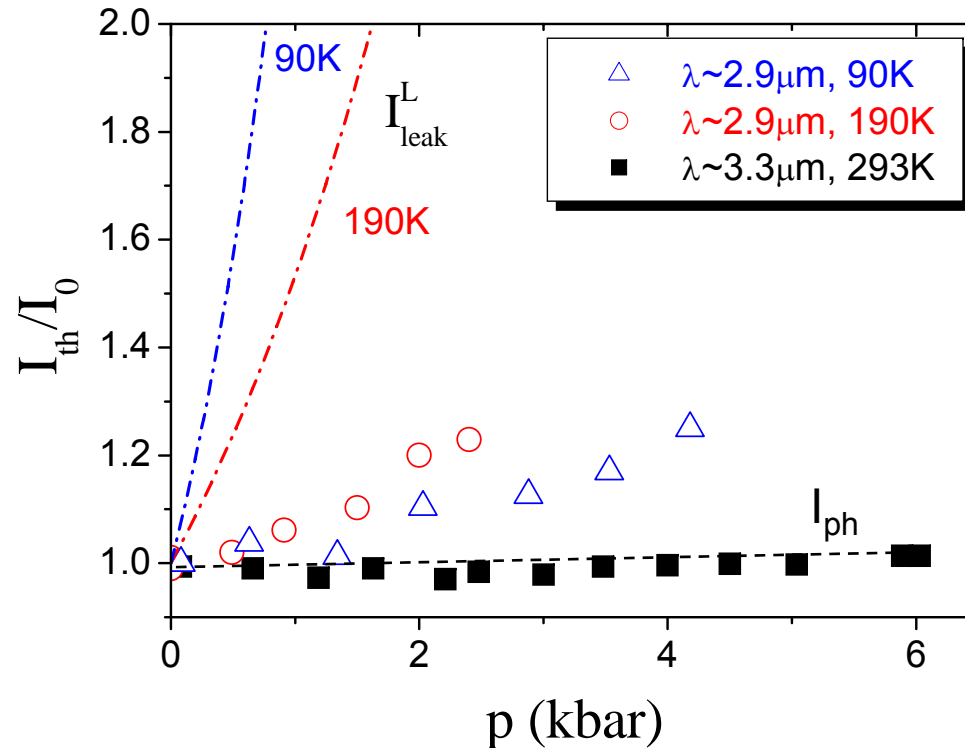
Indirect carrier leakage has approx.
 exponential pressure dependence
 with $dE_L - dE_\Gamma \approx -7\text{meV/kbar}$

Pressure dependence of I_{th}



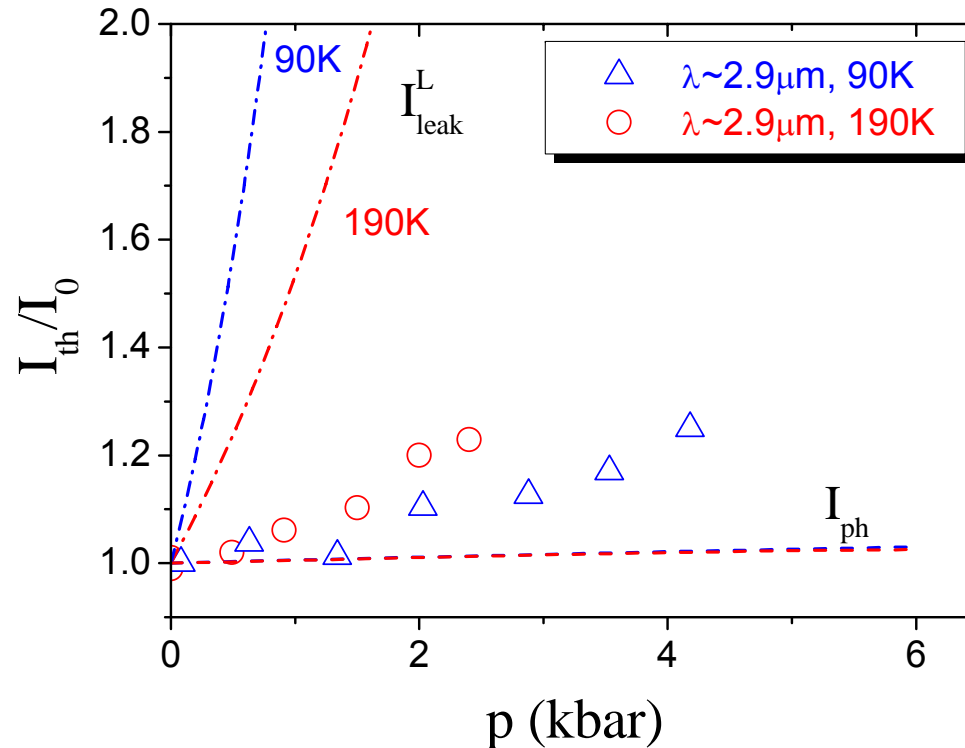
- I_{th} of 3.3 μm QCL at 293K is almost pressure independent even at RT indicating that I_{leak}^L is negligible in this device even at RT
- $I_{th} \sim I_{ph}$ - current due to intersubband carrier relaxation via LO phonon scattering with very weak pressure dependence

Pressure dependence of I_{th}



- I_{th} of 3.3 μm QCL at 293K is almost pressure independent even at RT indicating that I_{leak}^L is negligible in this device even at RT
- $I_{th} \sim I_{ph}$ - current due to intersubband carrier relaxation via LO phonon scattering with very weak pressure dependence
- In contrast, I_{th} in 2.9 μm QCL increases considerably with increasing pressure
- Laser operation stops at 2.4kbar at 190K and at 4.2kbar at 90K

Quantitative estimations of carrier leakage



- Dashed lines show calculated $I_{ph}(p)$ for 2.9 μm QCL at 90K and 190K

(similar calculations for GaAs/AlGaAs QCL - S. R. Jin et al, APL **89**, 221105, 2006)

- Estimate fractional contribution of leakage by assuming that:

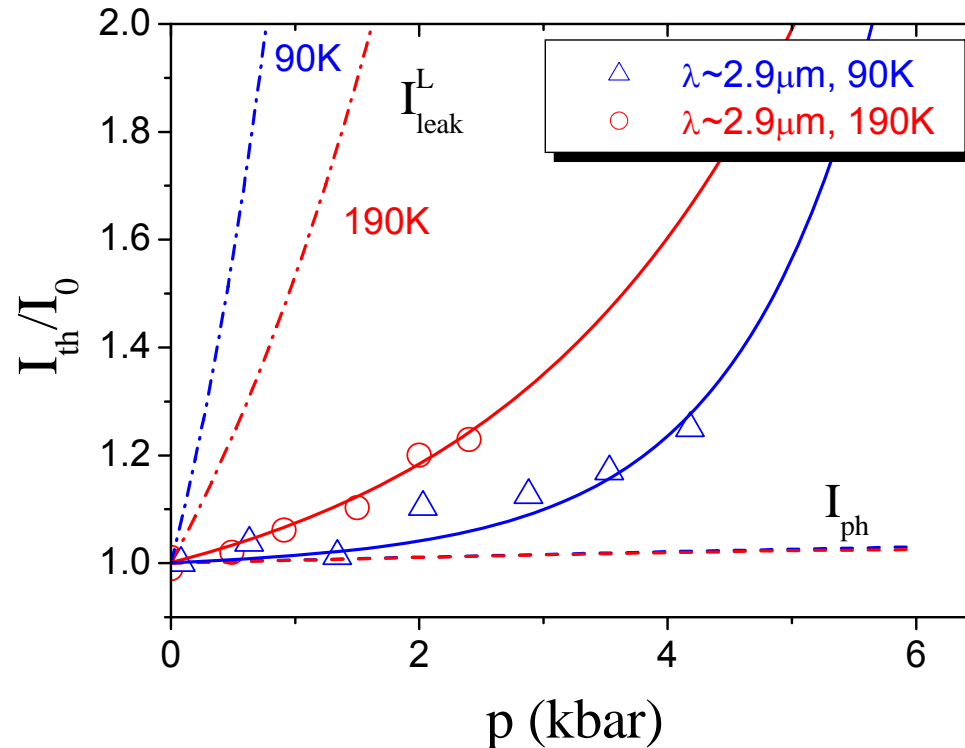
$$I_{th} = I_{ph} + I_{leak}^L$$

or

$$\frac{I_{th}}{I_0} = (1 - x) \cdot \frac{I_{ph}}{I_{ph0}} + x \cdot \frac{I_{leak}^L}{I_{leak0}^L}$$

x – fraction of I_{th} due to leakage at $p=0$

Pressure dependence of I_{th}

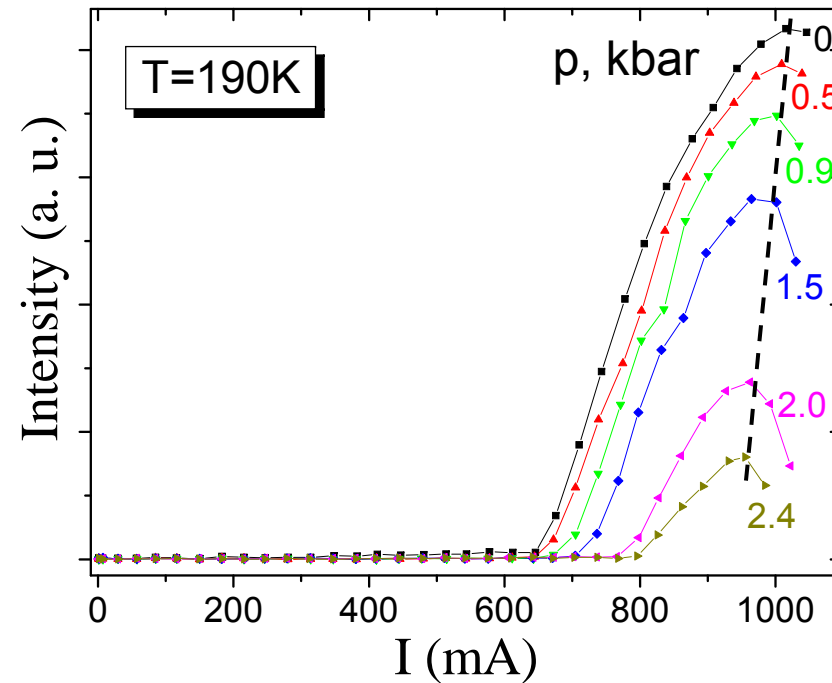
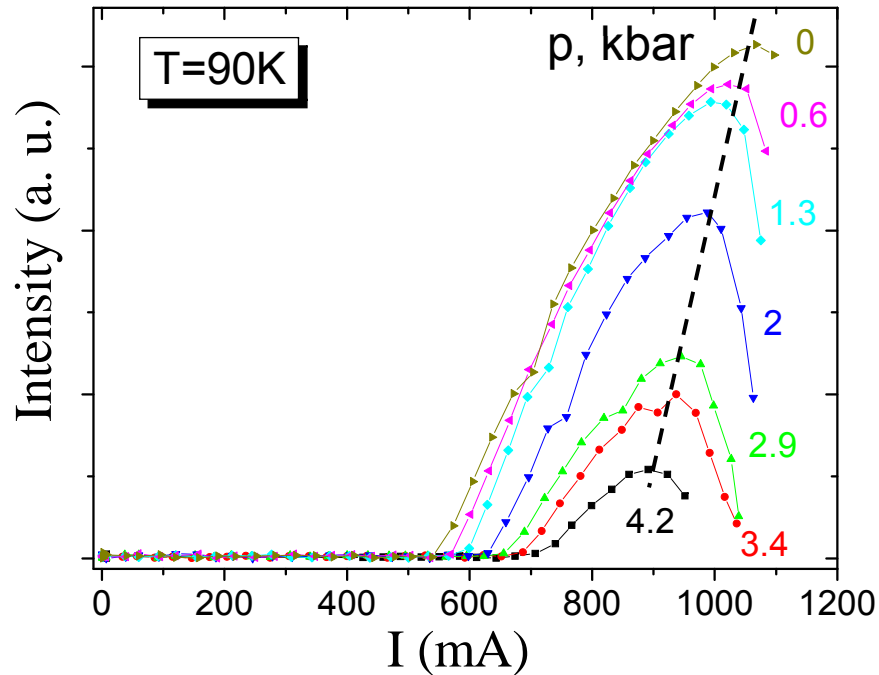


- $x = 0.006$ (0.6%) at 90K
- $x = 0.13$ (13%) at 190K
- By applying pressure we increase leakage current up to 29% of I_{th} at 190K and 20% of I_{th} at 90K

$$\frac{I_{th}}{I_0} = (1 - x) \cdot \frac{I_{ph}}{I_{ph0}} + x \cdot \frac{I_{leak}^L}{I_{leak0}^L}$$

x – fraction of I_{th} due to leakage at $p=0$

Pressure dependence of L-ls: 2.9 μm QCL



- Maximum available current, I_{max} , decreases with increasing pressure
- This is explained by the reduction of population of the upper level of lasing transition due to leakage into the L-valleys
- Such behaviour is not observed in the 3.3 μm QCLs

Conclusions

- The impact of Γ -L scattering on the temperature performance of short wavelength InAs/AlSb QCLs was investigated using high hydrostatic pressure.
- The results show evidence of that this loss mechanism is important in $\sim 2.9 \mu\text{m}$ lasers ($\sim 20\%$ of I_{th} at maximum T), but much less pronounced in $\sim 3.3 \mu\text{m}$ devices, enabling their superior temperature performance.
- We suggest that increased leakage into the L-minima is the reason for the even lower maximum operating temperature (140 K) of lasers emitting near $2.75 \mu\text{m}$.